



# Characterization of a Non-Thermal Plasma Source for the Use as a Mass Spec Calibration Tool and Non-Radioactive Aerosol Charger

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**Abstract.** In this study the charging efficiency of a radioactive and a non-radioactive plasma neutralizer (Gilbert Mark I plasma charger) have been investigated at various aerosol flow rates. The results were compared to classic theoretical approaches. In addition, the chemical composition and electrical mobilities of the charger ions have been examined using an atmospheric pressure interface - time-of-flight mass spectrometer (APi-TOF MS). A comparison of the different neutralization methods revealed an increased charging efficiency for negatively charged particles using the non-radioactive plasma charger with nitrogen as working gas. The mobility and mass spectrometric measurements show that the generated neutralizer ions are of the same mobilities and composition independent of the charging mechanism. It was the first time that the Gilbert Mark I plasma charger was characterized in comparison to the standard TSI X-Ray (TSI Inc, Model 3088) and a radioactive americium neutralizer. We observed that the plasma charger with nitrogen as working gas can enhance the charging probability for sub-10 nm particles. Consequently, the limit of detection of differential or scanning mobility particle sizers can be increased down towards to smaller sizes with the Gilbert Mark I plasma charger.

## 1 Introduction

Bipolar diffusion charging and neutralization is typically done by ionizing radiation, which exposes aerosol particles to high concentrations of positive and negative ions in the carrier gas (Jiang et al., 2014). Subsequent diffusion of the ions brings the aerosol to a stationary state charge distribution independent of their initial charge state (Cooper and Reist, 1973; Liu and Pui, 1974; Adachi et al., 1985; Reischl et al., 1996). If a high ion concentration and residence time is reached, a charge equilibrium inside the charger leads to a well known size-dependent charging probability (Fuchs et al., 1965). This stationary state charge distribution is of importance for the use of differential or scanning mobility particle spectrometers, which rely on accurate knowledge of the size-dependent charge fractions (Wang and Flagan, 1990; Jiang et al., 2014). Aerosol particles below 10 nm in diameter are typically difficult to neutralize and carry only one electrical charge at maximum (Wiedensohler, 1988). Quantitative particle detection in this size range is extremely challenging due to high diffusional losses. Hence, a higher charging efficiency is of importance to improve the signal intensity in the sub-10 nm regime.

The charging of small particles has also become a field of major interest in plasma physics. A vast number of studies about

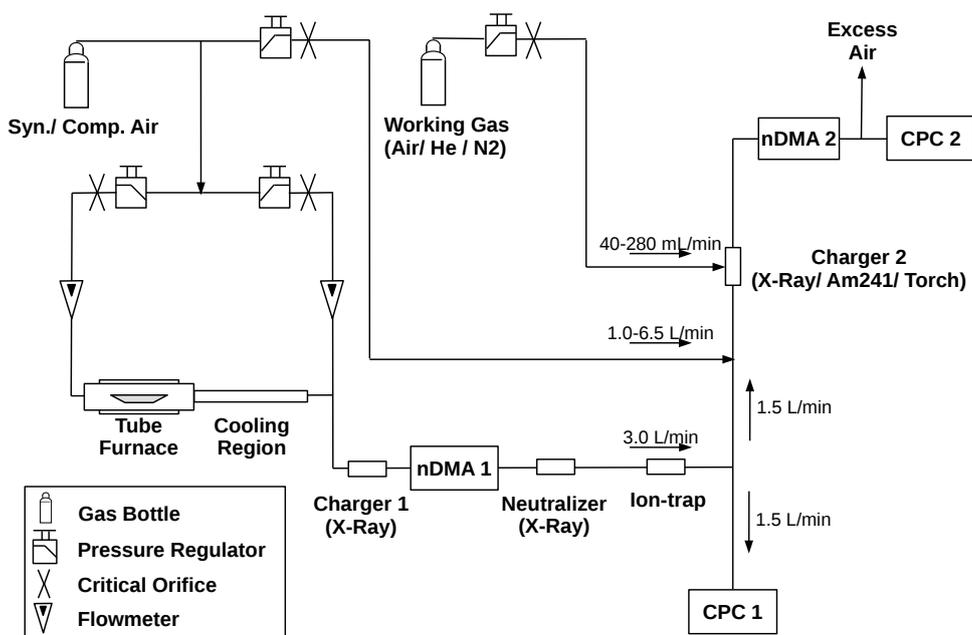


accumulating an electrical charge on dust particles have been published over the last years. Usually grain sizes ranging from  
25 some nm to several  $\mu\text{m}$  are considered in the experiments and theoretical models. Most of these works were focused on  
generating plasmas for industrial or space applications (Michau et al., 2016; Deka et al., 2017; Kopnin et al., 2018; Yaroshenko  
et al., 2018; Intra and Yawootti, 2019). However, one of the most recent developments is the application of plasma in aerosol  
related topics, such as plasma treatment of aerosol particles (Uner and Thimsen, 2017) and, as a novel topic, charging of  
aerosol particles or ionization of trace gas compounds (Spencer et al., 2015; Yang et al., 2016; Intra and Yawootti, 2019).  
30 Usually, corona-type discharges are used for the charging purposes in aerosol physics due to their capability of creating high  
charge densities even at atmospheric pressure. Furthermore, their reproducibility is very high and they are easy to construct and  
maintain. Low-temperature plasma ionization is known to cause little fragmentation and exhibits a low temperature increase to  
the surrounding (Harper et al., 2008). The non-thermal plasma in this work is produced by a high frequency generator which is  
separated by a dielectric barrier to the ground potential (Gruenwald et al., 2015). The discharge characteristics can be varied by  
35 changing the working gas of the Gilbert Mark I plasma charger (Gruenwald Laboratories GmbH). Plasma discharges in general  
are on-off devices that combine the simple handling of an X-ray charger with the achievable high ion density of a radioactive  
americium charger and even higher. Thus, an atmospheric plasma source is a well-suited device for the ionization process prior  
to aerosol size distribution measurements and mass spectrometric measurements.

In the past, various studies have characterized the charging probabilities and mobility spectrum of AC-corona, X-Ray or  
40 alpha-radiation based chargers (Wiedensohler et al., 1986; Steiner and Reischl, 2012; Kallinger et al., 2012; Kallinger and  
Szymanski, 2015). In this work, we present the experimental results obtained with the commercially available Gilbert Mark I  
plasma charger, which will be described in more detail in the next section. In addition, the chemical composition of charger  
ions of both polarities has been investigated and compared. Furthermore, the optical emission spectra were measured for the  
Gilbert Mark I plasma charger.

## 45 2 Experimental Setup

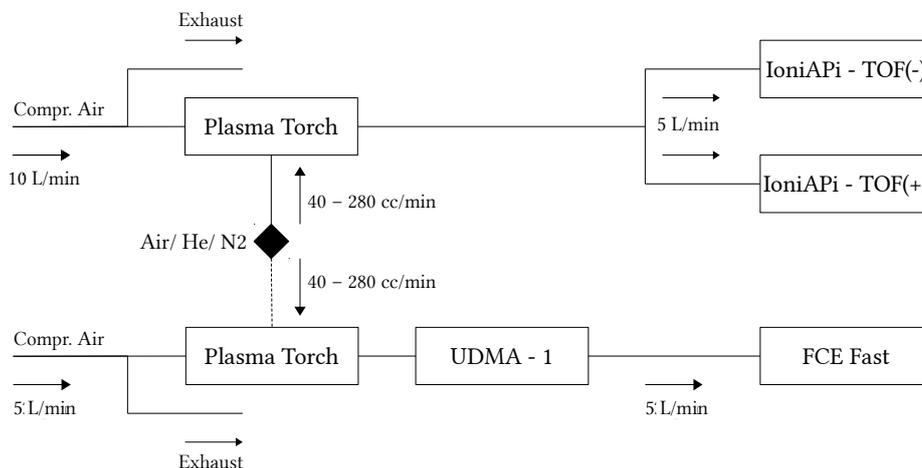
Here we report on size-dependent steady-state charging probability measurements of a non-thermal plasma source (Gilbert  
Mark I plasma charger, Gruenwald Laboratories GmbH, Austria), an americium 241 ( $^{241}\text{Am}$ ) charger and of a TSI Advanced  
Aerosol Neutralizer 3088 by means of a tandem DMA (Differential Mobility Analyzer) setup as depicted in Figure 1. Thereby,  
the charging efficiency of a standard TSI X-Ray neutralizer and an  $^{241}\text{Am}$  charger could be compared to the plasma source. The  
50 charging efficiency measurements were performed with sodium chloride and silver nanoparticles at various particle sizes and  
aerosol flowrates. The nanoparticles were generated with a tube furnace (Carbolite Gero GmbH & Co. KG, Germany) while  
synthetic air as well as dried and filtered compressed air were used as the carrier gas. An additional dilution flow allowed for  
controlling the particle concentration of the generated aerosol flow. Downstream of the aerosol generator the nanoparticles were  
charged with a TSI Advanced Aerosol Neutralizer 3088 and led to a nanoDMA which was operated as a classifier. A second  
55 TSI Advanced Aerosol Neutralizer 3088 neutralized the monodisperse aerosol particles after the nanoDMA. To secure that no  
charged particles remain in the aerosol flow an ion-trap was installed after the neutralizer. After the ion-trap the aerosol flow



**Figure 1.** Schematic of the experimental setup for the charging probability and particle size conservation measurements. See text for explanation.

was split in a way that 1.5 lpm were led to a CPC (TSI 3776 UCPC) which recorded the particle concentration at this point while the remaining aerosol flow was fed into the aerosol neutralizer under investigation. A dilution flow of synthetic air as well as compressed air allows for varying the flowrate before the aerosol flow gets to the neutralizer. The investigated neutralizers were  
60 switched in intervals of 10 minutes during the measurements to secure that the different devices were operated at comparable conditions. Afterwards the charged aerosol flow was led to a second nanoDMA, which was operated in scanning mode. A second CPC (TSI 3776 UCPC) recorded the particle concentration at this point of the setup. In addition to the switching of the different neutralizers, the working gas as well as the working gas flow of the atmospheric pressure plasma source were varied.

Complementary to the charging efficiency, the chemical composition of the charger ions were investigated by coupling the  
65 plasma torch with an API-TOF MS (Atmospheric Pressure Interface - Time of Flight Mass Spectrometer, Junninen et al. (2010); Leiminger et al. (2019), see upper panel in Figure 2), to analyze the chemical composition of the ions generated in positive and negative ion mode. Mobility spectra were recorded with a custom-built Faraday Cup Electrometer (FCE) (Winklmayr et al., 1991) with an improved response time of 0.1 s. By recording the ion spectrum with a Vienna-type high-resolution mobility analyzer (UDMA-1, Steiner et al. (2010)), the mobility equivalent diameter of the generated clusters could be analysed (see  
70 lower panel in Figure 2). Compressed filtered and dry air was used as carrier gas and the relative humidity (RH) was monitored



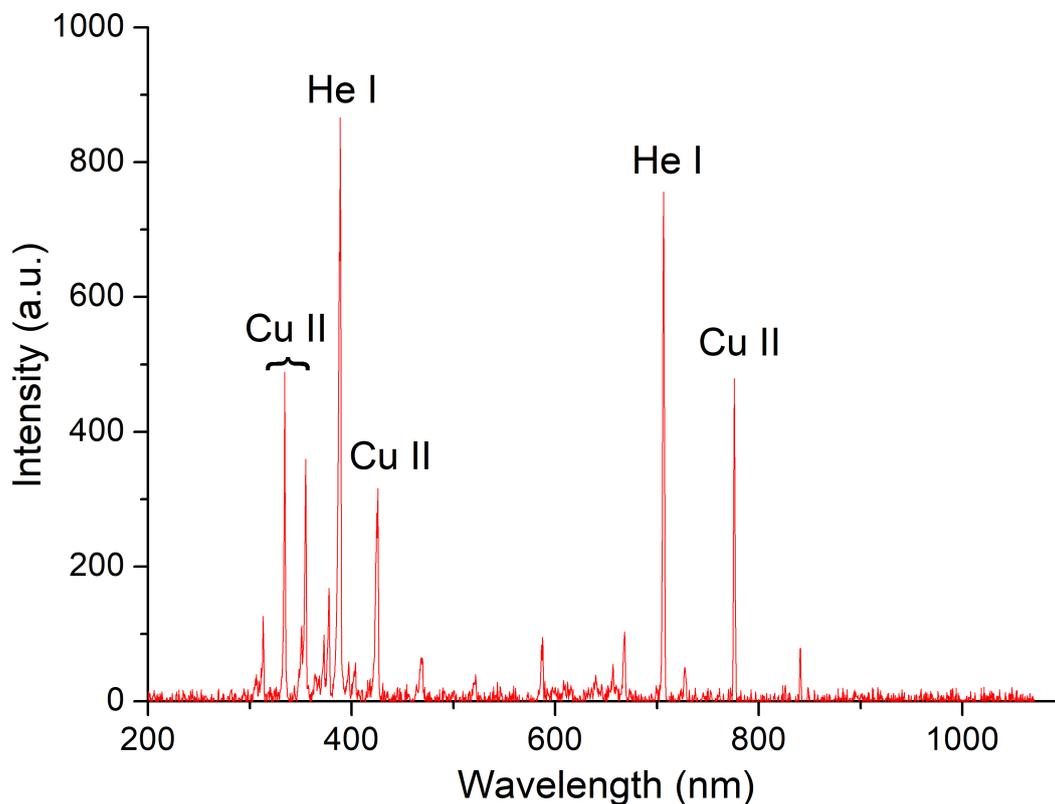
**Figure 2.** Schematic of the calibration setup for the plasma torch using an UDMA. The black rhombus marks the gas supply. Mass spectra of negative and positive ions were measured simultaneously using the IoniAPi-TOF in negative and positive ion mode at a flow rate of 10 L/min through the plasma torch (upper panel) for the three working gases and operational settings. The mobility spectra of the generated positive and negative ions were measured using the UDMA-1 (Steiner et al., 2010) coupled to a fast-response FCE at 5 L/min flow rate.

using SHT75 RH sensors with an accuracy of  $\pm 1.8\%$  and was kept below 2%. However, we can not exclude that in the closed loop sheath flow system of the UDMA-1 small amounts of water vapor remain. The calibration of the IoniAPi-TOF mass axis was performed using a bipolar electrospray source for the generation of tetra-heptyl ammonium bromide clusters (Fernández de la Mora and Barrios-Collado, 2017; Brilke et al., 2019). The UDMA-1 resolution power is 15 at the size of the THABr monomer, i.e. 1.45 nm mobility equivalent diameter (Flagan, 1999; Steiner et al., 2010). Due to the high ion concentration of a non-thermal plasma source a vast number of reactive species are created, especially when the plasma is ignited in air (Kurake et al., 2016). Most of these species will be ozone or nitrogen oxides because of the air's chemical composition. Hence, the optical emission spectra of the non-thermal plasma source were recorded, which is a non-invasive diagnostic technique that allows to gain insights into the composition of the plasma and the production of harmful gases like ozone and nitrogen oxides.

## 80 2.1 Results and Discussion

### 2.1.1 Optical Emission Spectroscopy

Optical emission spectroscopy (OES) was used to determine the ionization stages of the ions/molecules in the plasma. The measurements were performed with the HR2000+ES spectrometer from Ocean Optics. The light emission from the atmospheric pressure plasma source was collected with a 440  $\mu\text{m}$  fiber with a length of 2 m. The wavelength range of the spectrometer was between 200 and 1100 nm. Prior to the data acquisition the torch was switched on for about 30 seconds until no fluctuations in the spectra were visible. Each spectrum is the result of averaging over 50 scans to remove fluctuations. The plasma source itself was driven with a high frequent alternate current of 15 kHz with 825 V peak-to-peak voltage. The plasma was ignited



**Figure 3.** Typical OES spectrum in a helium plasma averaged over 50 scans. Experimental conditions: 15 kHz driving frequency, 825 volts peak-to-peak and 180 mL/Min He flow.

in helium of high purity (ALPHAGAZ 1 HELIUM,  $\geq 99.999\%$  (5.0), Air Liquide), which was fed into the plasma jet with a flow rate of 180 mL/min. The results for a typical OES spectrum for the described experimental conditions are depicted in  
90 Figure 3.

The most prominent emission lines were identified to be excited neutral helium (He I) and singly ionized copper (Cu II). The central wavelengths of the identified lines are listed in Table 1.

It can be concluded from the OES spectra that there is no ionization of aerosol particles facilitated by the carrier gas, since only neutral helium emission lines have been recorded. On the other hand, atoms from the copper high frequency antenna enter  
95 the plasma zone and are then ionized through electron impacts. The charged particles created from these processes (i.e. ions and secondary electrons), in turn, charge the aerosol particles in the gas stream. This can be explained by the large difference in ionization energies between copper (7.7 eV) and helium (24.6 eV). Since the electrons in non-thermal atmospheric pressure discharges have normally energies of just a few eV, the ionization of copper atoms is far more likely than of helium particles.



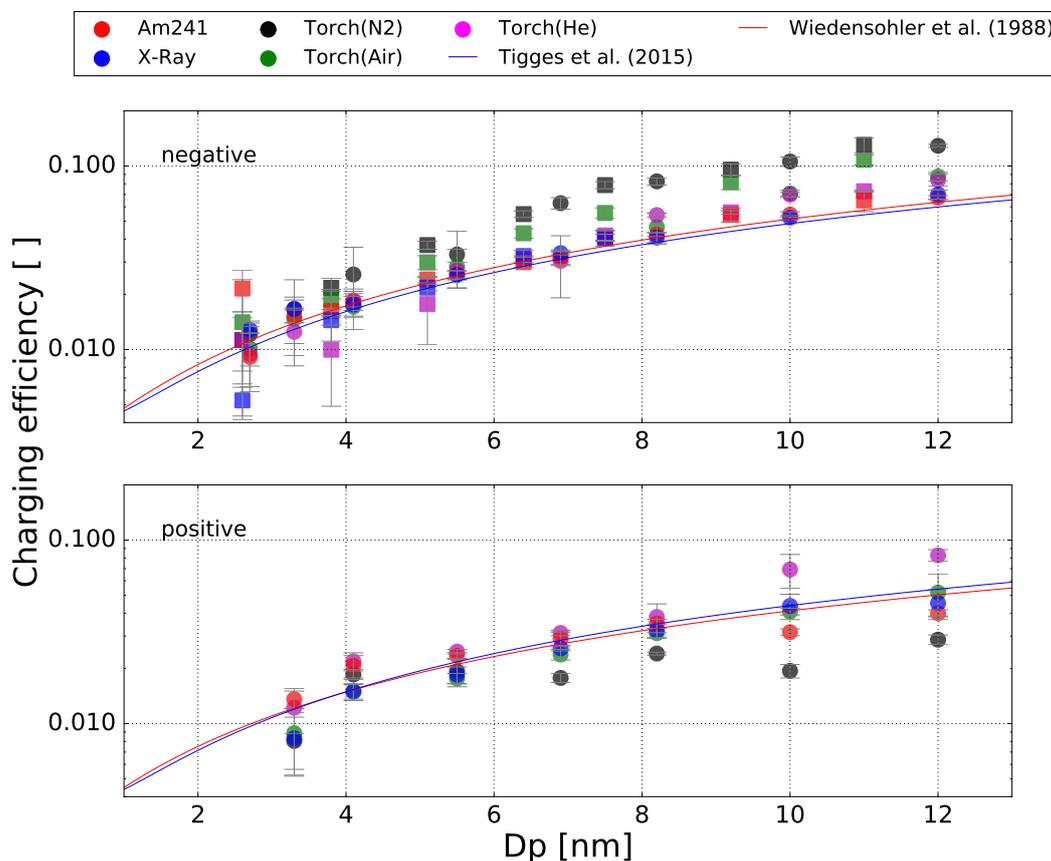
**Table 1.** Measured Central wavelengths compared to data from Kramida et al. (2013) and the associated particle species for the most prominent emission lines in Figure 3.

| Measured central wavelength [nm] | Wavelength from Kramida et al. (2013) [nm] | Particle Species |
|----------------------------------|--|------------------|
| 334.38                           | 334.4                                      | Cu II            |
| 354.86                           | 354.9                                      | Cu II            |
| 388.86                           | 388.9                                      | He I             |
| 425.67                           | 425.6                                      | Cu II            |
| 706.05                           | 706.5                                      | He I             |
| 755.56                           | 775.4                                      | Cu II            |

### 2.1.2 Charging Probability

100 The neutralizers were tested in the setup shown in Figure 1, in order to characterize the charging performance of the plasma torch (Gilbert Mark I plasma charger) and the  $^{241}\text{Am}$  as well as the soft X-Ray charger. The tandem DMA setup enabled us to charge a monodisperse aerosol and rotate the different neutralizers, which permits a direct comparison of the charging performance of the different devices. Two butanol-based CPCs (TSI 3776 UCPC) with reduced temperature settings to increase the particle counting efficiency were used (Tauber et al., 2019a). The particle number concentration was recorded before (CPC1, 105 See Fig. 1) and after charging (CPC2) by the tested charger (Charger 2). The charging efficiency was inferred from the ratio of the two CPCs (CPC2/CPC1) under consideration of the transmission and diffusional particle losses in the lines and DMA. In Tauber et al. (2019b) the particle counting efficiency of the CPCs used here was determined, and the results obtained were corrected for the CPC detection efficiency. The different neutralizers were tested with positively and negatively charged Ag and NaCl particles of different particle sizes in the sub-10 nm regime. In addition, the plasma torch was operated with different 110 working gases. The results of the charging efficiency when using the three different working gases are displayed in Figure 4. Theoretical charging probabilities from Tigges et al. (2015) and Wiedensohler (1988) were added here.

Concerning negatively charged particles with a diameter between 4-12 nm the  $^{241}\text{Am}$  and the soft X-Ray charger are in agreement with the theoretical curves especially with Wiedensohler (1988). The plasma torch on the other hand achieves higher charging efficiencies in this regime with big differences between the used working gases. With helium the charging 115 efficiencies were higher than the theoretical values, especially for particle sizes bigger than 8 nm. Ambient air yields even higher charging probabilities but displays a significant dependency on the particle type as the values for NaCl particles are a lot higher than for Ag particles. The highest charging efficiencies for negatively charged particles were achieved with nitrogen as working gas. With nitrogen the plasma torch achieved charging efficiencies that were up to 50 % higher than the common devices and showed no dependency on the particle type. As the particle concentrations for very small diameters between 2-4 120 nm are below 5000 #/ccm, the charging efficiencies of the different devices vary strongly at these particle sizes. However,



**Figure 4.** Measured charging efficiencies for the different aerosol chargers for negatively and positively charged Ag particles (dots) and negatively charged NaCl particles (squares) with mobility diameters less than 12 nm.

the associated data points are in agreement with theory. This is not the case for NaCl particles as one can clearly see big variations in the charging efficiencies of the different charging devices. Those asymmetries almost certainly are owed to the low concentrations of NaCl particles with these sizes.

The charging efficiencies of the plasma torch for positively charged particles strongly differ from those of negatively charged particles. Again, the measured charging efficiencies of the soft X-Ray charger match almost perfectly with the theory for diameters between 4-12 nm. The  $^{241}\text{Am}$  charger also agrees with the theoretically predicted values for diameters between 6-8 nm but shows asymmetries for larger particle sizes. The charging efficiency for positively charged particles decreases at larger sizes and results in lower values compared to theory for the  $^{241}\text{Am}$  charger. The charging efficiencies of the plasma torch strongly depend on the type of working gas for positively charged particles. Evidently, the data using helium agrees well in the size regime between 4 and 10 nm. For diameters between 10-12 nm the charging efficiencies even exceed the predicted values. The data suggest that no dependency on the charger ion polarity is present, when helium is used as working gas in the plasma torch. The data for compressed air match with the theoretical curves for the whole size range, and therefore behave similar

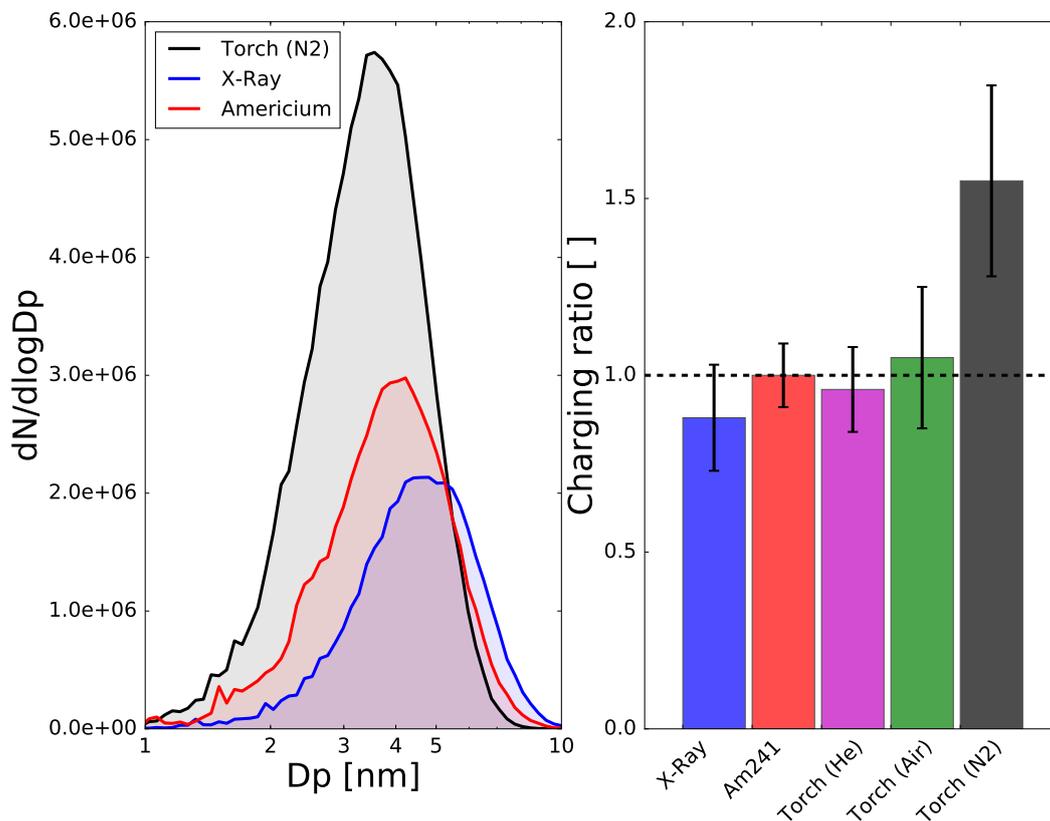


to the soft X-Ray charger. As Figure 4 clearly shows, this will change drastically if nitrogen is used as the working gas. The charging efficiencies in positive polarity with nitrogen are significantly lower than with other chargers and other working gases in the whole size regime.

Wiedensohler and Fissan (1991) have shown that the predicted charging probabilities of NaCl and Ag particles strongly depend on the used carrier gas and the ion mass. During the ionization process positive ions and free electrons are formed from molecules in the carrier gas and ionized copper atoms. These primary ions attach to other molecules, as for example  $H_2O$ ,  $CO_2$ , oxygen and halogens, and form bigger ion clusters that afterwards stick to the investigated particles. Wiedensohler and Fissan (1991) have shown that the variation of the ion masses leads to different theoretically predicted charging probabilities. For nitrogen as carrier gas they discovered a large dependency of the charging probability on the ion masses. Similar asymmetries are observed in our data when comparing the measured charging efficiencies of the plasma torch with nitrogen as working gas to the  $N_2$  results of Wiedensohler and Fissan (1991). As Figure 3 shows, the plasma torch forms copper ions and free electrons which charge aerosol particles in the carrier gas. The significantly different masses of these ions may account for the differing charging efficiencies that are accomplished with nitrogen for the negatively as well as for the positively charged particles according to Wiedensohler (1988). As the working gas flow is exposed to a high frequency electrical field before it mixes with the aerosol flow, the ions can form in a pure nitrogen environment almost like in the mentioned case of Wiedensohler and Fissan (1991). A similar but smaller effect was observed in atmospheric air as carrier gas. This mechanism would also explain the better charging efficiencies with ambient air as working gas.

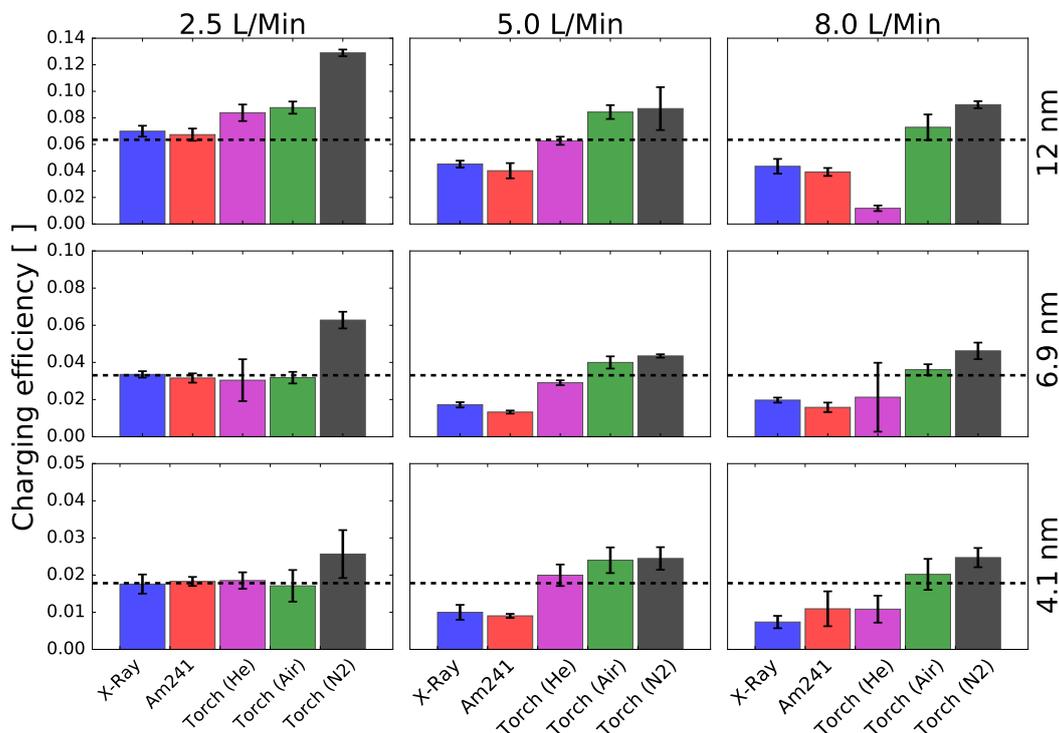
Figure 5 depicts the recorded aerosol size distribution averaged over numerous measurements, as well as the charging efficiency of the individual chargers normalized to the  $^{241}\text{Am}$  charger recordings. The inversion of the size distribution data was made according to Petters (2018). These diagrams permit a descriptive comparison of the different charging devices. The left plot reveals a shift of about 1 nm in the maximum of the recorded size distribution between the X-Ray and the plasma charger. In addition, a significantly higher peak with the plasma torch was recorded which is most likely due to the higher charging efficiency at smaller sizes. Since the neutralizers were rotated periodically in multiple cycles, the possibility of systematic uncertainties in the actual size distribution was minimized. Therefore, the raised signal of the plasma torch can be attributed to a generally higher charging efficiency. The right plot simplifies a direct comparison of the plasma torch operated with different working gases and the soft X-Ray as well as the  $^{241}\text{Am}$  charger as the calculated charging efficiencies are normalized to the  $^{241}\text{Am}$  charger. The X-Ray charger is slightly less efficient than the radioactive  $^{241}\text{Am}$  charger but still within the uncertainty range. This decrease can be attributed to the performance reduction during continuous operation which typically occurs during long measurement cycles. Compared to conventional chargers the plasma torch proves to charge slightly better with air as the working gas and less efficient with helium. With nitrogen as working gas, the plasma torch charges up to 50 % more than with air and helium and more than the other tested neutralizers.

Figure 6 shows the dependence of the charging efficiency for different aerosol flowrates and particle sizes. The charging efficiencies for particles in the sub-50 nm regime show a significant dependence on the aerosol flowrate. The different neutralizers have proven to be more sensitive to varying flow rates and a reduced residence time in the ionizing atmosphere (He and Dhaniyala, 2014; Kallinger et al., 2012; Kallinger and Szymanski, 2015).



**Figure 5.** The plot on the left side represents the same recorded Ag aerosol size distribution with a UCPC, depending on the mobility diameter for the plasma torch (N<sub>2</sub>), X-Ray and americium neutralizers. On the right-hand side all conducted measurements with the same aerosol distribution are normalized to the <sup>241</sup>Am results for a direct comparison.

Figure 6 reveals that the aerosol charging is most efficient for a flowrate of 2.5 L/Min and the charging efficiency decreases for higher flowrates. Kallinger and Szymanski (2015) and Jiang et al. (2014) also measured the flowrate dependence of different  
170 chargers. In the study by Kallinger and Szymanski (2015) an increased charging efficiency for the americium neutralizer for a flowrate of 5.0 L/Min was found. In addition, contradictory to our findings, a reduced charging performance for the X-Ray  
neutralizer was not observed by Kallinger and Szymanski (2015) and Jiang et al. (2014). A reason for that could be a reduced charging due to low power output of the X-Ray tube. Since the last repair and calibration of the analyzed TSI 3088 neutralizer the operating run-time was about 371 hours. Furthermore, in this work, we used a different aerosol generation method  
175 compared to the above mentioned studies which could lead to different charging mechanism. Especially in the sub-10 nm size range different chemo-physical interactions might lead to unforeseen results which should be investigated in future studies. The plasma torch also achieves the highest charging efficiencies for lower flowrates but seems to be not as sensitive to the aerosol flow compared to the other devices. Especially air and nitrogen have proven to be the most robust options as working gases.



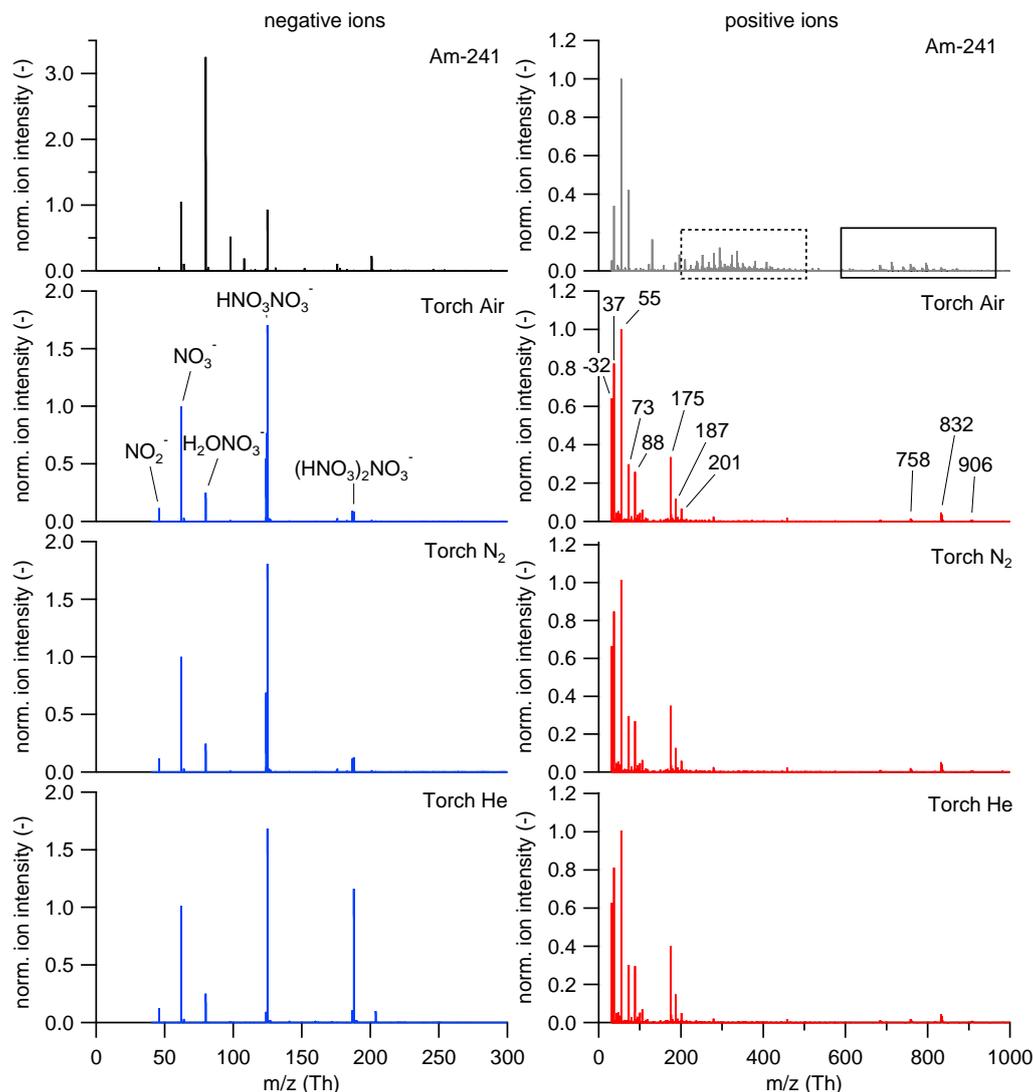
**Figure 6.** Measured charging efficiencies for the different aerosol chargers and aerosol flow rates for negatively charged Ag particles with mobility diameters of 12, 6.9 and 4.1 nm. The dotted black line represents the theoretical charging efficiency according to Wiedensohler (1988) for the 3 different mobility diameters.

### 2.1.3 Charger ion chemical composition

180 The ion properties of ionic molecular clusters produced in the plasma torch were investigated by means of electrical mobility and mass spectrometry. Mobility spectra and mass spectra were recorded for positive and negative ions and compared to the resulting spectra from ions produced in the  $^{241}\text{Am}$  charger. Figure 7 shows the mass spectra for negative (left) and positive (right) ions generated by the  $^{241}\text{Am}$  charger (first panel) and the plasma torch for the three different working gases (second to fourth panel) using the setups shown in Figure 2 at the working gas flow settings presented in Table S1 and S2 in the supporting information (SI). The positive mass spectra were normalized to the nitrate ion ( $\text{NO}_3^-$ ) peak at an integer mass of 62 Th and the negative mass spectra to the  $(\text{H}_2\text{O})_2 \cdot \text{H}_3\text{O}^+$  water cluster at an integer mass of 55 Th. The negative mass spectra are dominated by the nitrate ion,  $\text{NO}_3^-$ , and its dimer, trimer and water cluster (see labels in second panel in Figure 7). The three spectra for charger ions produced from the torch exhibit the same major peaks with the nitrate ion trimer peak being highest when He is used as working gas (see fourth panel). This observation may be a result of the different operational settings of the plasma torch when He is used as working gas (see Table S1 and S2). The negative ion spectrum of the Am241 charger reveals the same major peaks as the plasma torch negative mass spectra. Similar results have been found by a study investigating

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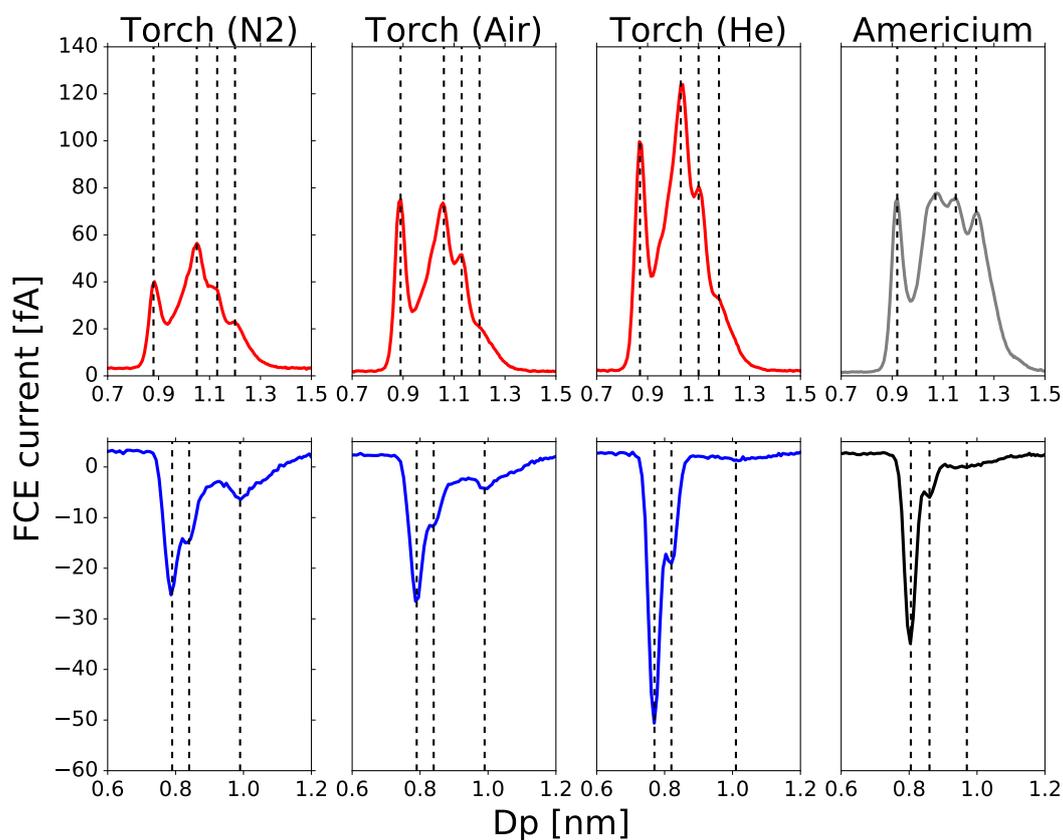
**Figure 7.** Negative (left) and positive (right) mass spectra of ions generated by the <sup>241</sup>Am charger (first panel) and the plasma source for three different working gases, synthetic air (second panel), N<sub>2</sub> (third panel) and He (fourth panel) were measured using the setup in the upper panel of Figure 2. The mass spectra were averaged over 1 hour each. The identified compounds are labeled in the second panel and are presented in Table 2. The negative mass spectrum was normalized to the NO<sub>3</sub><sup>-</sup> ion (integer mass 62 Th), the positive mass spectrum was normalized to the (H<sub>2</sub>O)<sub>2</sub>·H<sub>3</sub>O<sup>+</sup> cluster (integer mass 55 Th). The H<sub>3</sub>O<sup>+</sup> ion is not displayed here since it was not covered by the set mass range of the ioniAPI-TOF. The dashed square box marks unidentified masses in the positive <sup>241</sup>Am mass spectrum and the solid square box shows the silicone compounds that are listed in Table 2.

the chemical composition of ions produced by a corona discharge (Manninen et al., 2011). The identified major peaks of the positive mass spectra are listed in Table 2. In the lower mass range between 40 - 80 Th, protonated water, H<sub>3</sub>O<sup>+</sup>, and



195 water clusters thereof dominate the spectrum for the  $^{241}\text{Am}$  charger and the three spectra of the plasma torch. The elemental  
composition of four major peaks in the positive spectra of the plasma torch at integer masses of 88, 175, 187 and 201 Th were  
200 identified as carbonaceous compounds. In the higher mass range between 600 - 1000 Th, a set of major peaks was identified  
as silicone compounds. The same peaks were identified in the positive mass spectra of ions produced from the  $^{241}\text{Am}$  charger  
(solid black box in Figure 7). This observation was also made by Manninen et al. (2011), who explain these peaks as a result of  
contamination from silicone tubing. Silicone tubing is oftentimes used in aerosol measurements and can cause artifacts because  
of degassing of siloxanes (Asbach et al., 2016). However, a range of peaks in the mass window from approximately 200 to 500  
Th remains unidentified (dashed black box). These carbonaceous compounds have a positive mass defect and likely arise from  
ionization of constituents in the pressurized air that was used as carrier gas.

205 The chemical composition of the plasma-generated ions was found to be independent of the choice of the working gas as  
shown in Figure 7. Also, the averaged electrical mobility measurements (averaged over 10 scans) conducted with the experi-  
mental setup depicted in Figure 2 (lower panel) revealed identical peaks. The mobility spectra for the plasma torch using  $\text{N}_2$ ,  
He and air as working gas and the  $^{241}\text{Am}$  charger are presented in Figure 8. Similar results for different bipolar charging de-  
vices have been found by Kallinger et al. (2012) and Steiner and Reischl (2012). The latter have analysed the effects of carrier  
gas contaminants on the charging probability which influences the electrical mobility spectrum. One of the analysed TSI X-Ray  
charger showed a different mobility spectrum compared to the other analysed neutralizers. According to our mass spec anal-  
210 ysis, this is due to ammonium sulfate contamination's from previous experiments which results in masses from  $(\text{NH}_4)_2\text{SO}_4^-$   
to  $(\text{NH}_4)_{14}\text{SO}_4^-$ .



**Figure 8.** Mobility distributions of the charger ions generated by the plasma torch using N<sub>2</sub>, He and air as working gas and the <sup>241</sup>Am charger. The upper panel shows the mobility diameter distribution of the positive charger ions, the lower panel presents the mobility diameter distribution of the negative charger ions.



**Table 2.** Overview of major negative and positive compounds in the mass spectra recorded using the ioniAPI-TOF in positive and negative ion mode.

| Negative Ions    |                          | Positive Ions    |                                |
|------------------|--------------------------|------------------|--------------------------------|
| Integer m/z (Th) | Molecular Formula        | Integer m/z (Th) | Molecular Formula              |
| 46               | $NO_2^-$                 | 32               | $O_2^+$                        |
| 62               | $NO_3^-$                 | 37               | $H_2O \cdot H_3O^+$            |
| 80               | $H_2O \cdot NO_3^-$      | 55               | $(H_2O)_2 \cdot H_3O^+$        |
| 124              | -                        | 73               | $(H_2O)_3 \cdot H_3O^+$        |
| 125              | $HNO_3 \cdot NO_3^-$     | 88               | $C_4H_{10}NO^+$                |
| 187              | -                        | 175              | $C_4H_9NO \cdot C_4H_{10}NO^+$ |
| 188              | $(HNO_3)_2 \cdot NO_3^-$ | 187              | $C_{10}H_{21}NO_2^+$           |
|                  |                          | 201              | $C_{11}H_{23}NO_2^+$           |
|                  |                          | 610              | $(SiOC_2H_6)_8OH_2^+$          |
|                  |                          | 684              | $(SiOC_2H_6)_9OH_2^+$          |
|                  |                          | 758              | $(SiOC_2H_6)_{10}OH_2^+$       |
|                  |                          | 832              | $(SiOC_2H_6)_{11}OH_2^+$       |
|                  |                          | 906              | $(SiOC_2H_6)_{12}OH_2^+$       |



### 3 Conclusions

The presented measurements conducted with a non-thermal plasma source, americium and X-Ray neutralizer have shown that the different charging mechanisms lead to the same ion species. According to the API-MS measurements the same chemical compounds were formed with the different charger mechanisms. At this point it should be noted that the chemical composition of the charger ions is affected by the tubing material and contaminations as discussed by Steiner and Reischl (2012).  
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It should be noted that the analyzed chemical composition of the neutralizer ions did not lead to changes in chemical composition even with increased ozone concentration caused by the plasma torch. By switching the working gas to nitrogen an increased charging efficiency could be recorded for negatively charged particles. In accordance with Mathon et al. (2017), the ozone concentration can be reduced to ambient conditions with nitrogen which is beneficial for the commercial use. For future studies, the influence of the variation of the operational settings and the resulting ozone concentrations of the plasma torch remains to be investigated.  
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The higher charging probabilities for negatively charged particles can be attributed to differences in the electrical mobility. According to Wiedensohler (1988) the ratio of ion mobilities, given by the ion mobility of positively charged particles divided by the ion mobilities of negatively charged particles, yields a constant value of 0.875. Consequently, the ion mobilities for negatively charged particles are on average higher. The increased charging efficiency for negatively charged particles and decreased charging efficiency for positively charged particles was also measured by Wiedensohler and Fissan (1991) for Kr85 bipolar neutralizer in nitrogen. According to our measurements, similar results could be found for the plasma torch with nitrogen as working gas. The charging efficiency of the non-thermal atmospheric plasma source indicated a weak aerosol flow dependence when operated with nitrogen or compressed air in comparison to the americium and X-Ray neutralizers.  
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In summary, with different experimental approaches we were able to quantitatively characterize the Gilbert Mark I plasma source with nitrogen, helium and air as working gas. In addition, a commonly used X-Ray neutralizer and a radioactive americium neutralizer were analyzed for comparison. The highest charging efficiencies for negatively charged particles were found for the Gilbert Mark I plasma charger with nitrogen as working gas. Our results also reveal the importance of well characterized and clean neutralizers to avoid any misinterpretation of experimental data especially in the sub-10 nm size range.  
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*Data availability.* Supplementary data associated with this article can be found in the online version.

*Author contributions.* Christian Tauber designed the setup, Christian Tauber and David Schmoll performed the charging efficiency experiments, Johannes Gruenwald and David Schmoll performed the OES measurements, Christian Tauber and Sophia Brilke performed the mobility distribution measurements, Sophia Brilke and Daniela Wimmer performed the chemical composition measurements, Christian Tauber, David Schmoll, Johannes Gruenwald, Sophia Brilke, Peter Josef Wlasits, Daniela Wimmer and Paul Martin Winkler were involved in the scientific interpretation and discussion, and Christian Tauber, David Schmoll, Johannes Gruenwald, Sophia Brilke, Peter Josef Wlasits, Daniela Wimmer and Paul Martin Winkler wrote the manuscript.  
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*Competing interests.* The authors declare that they have no conflict of interest.

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