

Abstract

This paper proposes a method to obtain spectral sky radiances, at three wavelengths (464, 534 and 626 nm), from hemispherical sky images. Images are registered with an All-Sky Imager installed at the Andalusian Center for Environmental Research (CEAMA) in Granada (Spain). The methodology followed in this work for the absolute calibration in radiance of this instrument is based on the comparison of its output measurements with modelled sky radiances derived from the Libradtran/UVSPEC radiative transfer code under cloud-free conditions. Previously, in order to check the goodness of the simulated radiances, these are compared with experimental values recorded by a CIMEL sunphotometer. In general, modelled radiances are in agreement with experimental data, showing mean differences lower than 15 % except for the pixels located next to the sun position that show larger errors.

The comparison between the output signal of the All-Sky Imager and the modelled sky radiances provides a calibration matrix for each image. The variability of the matrix coefficients is analyzed, showing no significant changes along a period of 5 months. Therefore, a unique calibration matrix per channel is obtained for all selected images (a total of 705 images per channel). Camera radiances are compared with CIMEL radiances, finding mean absolute differences between 2 % and 15 % except for pixels near to the Sun and high zenith angles. We apply these calibration matrices to three images in order to study the sky radiance distributions for three different sky conditions: cloudless, overcast and partially cloudy. Horizon brightening under cloudless conditions has been observed together with the enhancement effect of individual clouds on sky radiance.

AMTD

5, 1873–1905, 2012

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1 Introduction

Solar radiation is the most relevant term in the energy balance of the Earth-atmosphere system, being responsible of, e.g. photosynthesis in plants and vitamin D synthesis in humans (Webb, 2006). Nevertheless, it is well known that the ultraviolet range of solar radiation produces detrimental biological effects on human beings, animals and plants (Diffey, 1991). Solar radiation reaches the top of atmosphere in a fixed direction (direct radiation) however, while it is crossing the atmosphere, this radiation is diffused in other directions (diffuse radiation) due to scattering processes with gases, aerosol and clouds (Lenoble, 1993). Knowledge and measurements of the angular distribution of sky radiance (skylight at solar wavelengths, i.e. scattered sunlight) are important since the shape of human beings, animals, plants, etc. is not regular or oriented to the Sun. Therefore the sky radiance distribution has an outstanding role in the evaluation of radiation reaching complex targets like the human body or studies directed towards the development of solar energy systems.

The sky radiance depends on aerosol optical properties and thus these properties can be derived from sky radiance measurements. In this sense, different authors (e.g. Nakajima et al., 1996; Dubovik and King, 2000) proposed an inversion algorithm for the retrieval of aerosol optical properties using sky radiance measurements in the almucantar (the circle parallel to the horizon with the zenith angle equal to the solar zenith angle, SZA). Olmo et al. (2008b) developed an inversion method to obtain the columnar aerosol size distribution, the single scattering albedo (SSA), the phase function (PF) and the asymmetry parameter from sky radiance measurements in the principal plane (the plane perpendicular to the horizon that crosses the SZA and zenith) using non-spherical particles approximation. Therefore, the sky radiance measurements are also useful to retrieve the aerosol optical properties.

Clouds produce the strongest changes in solar radiation with a substantial decrease of the direct component and an increase of the diffuse radiation, which lead to a decrease in global radiation (Alados et al., 2000; Alados-Arboledas et al., 2003; de Miguel

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et al., 2011b). The influence of the cloudiness on solar radiation has a significant spectral dependence, being weaker when the wavelength is shorter (Foyo-Moreno et al., 2001; Bilbao et al., 2011; de Miguel et al., 2011a). On the other hand, the clouds can also produce the enhancement of the solar radiation at surface, reaching levels even higher than its value at the top of atmosphere (Piacentini et al., 2010; Antón et al., 2011a). Therefore, studies about the behaviour of sky radiance and clouds are important for a better understanding of cloud-radiation interaction.

Several authors have analyzed the sky radiance under different sky conditions using several types of instruments. For instance, Grant and Heisler (1997) and Grant et al. (1997a, b) used silicon photodiodes with filters, mounted on a platform that provided for rotational and inclinational motion, to take broadband radiance measurements in the ultraviolet range under obscured overcast, translucent overcast, and clear skies. Their system took a complete measurement of sky distribution (a grid of 10° zenith and 3° in azimuth) in half hour. Vida et al. (1999) tested a cloudless radiance model using radiance measurements taken by a pyroelectric radiometer (modified to present a 5° effective half angle field of view) and they found the highest differences between modelled and experimental radiances close to the Sun. Weihs et al. (2000) took spectral radiance measurements under different sky conditions using a spectroradiometer connected to a tube of foreoptics (1° field of view) mounted in a tracker programmed to cover a full sky radiance distribution (a grid of 10° in zenith and azimuth) in 15 min. Wuttke and Seckmeyer (2006) recorded measurements of sky luminance and spectral radiance using a Skyscanner that took 150 points, in 40 s, evenly distributed across the sky. Both works (Weihs et al., 2000; Wuttke and Seckmeyer, 2006) compared radiance measurements with radiative transfer models based on the discrete ordinate radiative transfer algorithm, DISORT (Stamnes et al., 2000), and the results indicated that circum-solar region shows the highest errors, and the input parameters would have to be known with greater accuracy (e.g. they do not use PF as input).

Sky cameras, or sky imagers, are devices that combine a digital camera with a fish eye lens, or a hemispherical mirror in some cases, that takes pictures of the full

is presented in Sect. 4. Results and conclusions are in Sect. 5, where the obtained calibration matrices are tested with measurements and a qualitative analysis of three sky images is presented. Finally, the conclusions are summarized in Sect. 6.

2 Site, instrumentation and data

2.1 Site

The measurements used in this paper were taken by the Atmospheric Physics Group (GFAT, <http://atmosfera.ugr.es>) of the University of Granada on the rooftop of the Andalusian Center for Environmental Studies (CEAMA, 37.17° N, 3.61° W, 680 m a.s.l.). Experimental site and its local conditions were described in detail by Lyamani et al. (2010, 2011).

2.2 Instrumentation

2.2.1 The All-Sky Imager

The All-Sky Imager designed and deployed by the GFAT is operating at our experimental site from 2005 and basically consists of a digital color CCD camera with a fish eye lens encapsulated in an environmental housing that is temperature regulated with a peltier cell. The system is installed in a Sun tracker (2AP model from Kipp & Zonen) to block the direct solar radiation using a shadow ball. A transparent acrylic dome protects the sky imager from the weather conditions like rain or snow. This dome was changed by a glass dome (non-acrylic) in December 2010 due to degradation problems, being the new dome more stable. The camera body is a RETIGA 1300C (QImaging) has a CCD sensor with three channels: Red, Green and Blue. The CCD sensor is the model ICX085AK from SONY working with a filter that blocks the wavelengths in the infrared region (standard configuration of the camera). The spectral responses of the camera channels are shown in Fig. 1, being centred around 450 nm, 550 nm and 650 nm with

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size distribution, complex refractive index, PF and SSA at different wavelengths following the AERONET procedures (Holben et al., 1998; Dubovik and King, 2000).

Calibration of this instrument was performed annually by AERONET-RIMA network. More details about CIMEL CE-318-4 can be found in Holben et al. (1998).

5 2.3 Data

Different measurements were used in this work. From the measurements recorded by the All-sky Imager, we only work with those recorded during the year 2011 in order to guarantee images with the glass dome.

10 On the other hand, the aerosol information was obtained from the CIMEL sunphotometer. Thus, the Angstrom's α parameter was calculated using AOD at 440 and 675 nm, and the Angstrom's β parameter was obtained taking into account α and AOD at 500 nm:

$$\alpha = \frac{\log \left(\frac{\text{AOD}_{440}}{\text{AOD}_{675}} \right)}{\log \left(\frac{675}{440} \right)}, \quad (1)$$

$$\beta = \frac{\text{AOD}_{500}}{0.5^{-\alpha}}. \quad (2)$$

15 In addition, the sunphotometer sky radiances measured with the 440 nm and 675 nm filters were provided by AERONET, but the products of the network are the radiances at the nominal wavelengths of 441 nm and 677 nm due to the filter bandwidth. Radiance measurements using the 500 nm filter (nominal wavelength of 501 nm) are only available along two months in 2011, but they were used to test the calibration method
20 in Sect. 5.1. Additionally, SSA values were also used in the work and, how they are retrieved using sky radiances, AERONET provides these values at the nominal wavelengths of 441 and 677 nm. The SSA values are only available at 441 and 677 nm, but we used the SSA at 550 nm, which was estimated as the mean between SSA441

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and SSA677. Similar method was followed for the PF values, which are only available at 441 and 677 nm (nominal wavelengths) in AERONET, and thus, PF550 was estimated as the mean between PF441 and PF677 for each scattering angle. Finally, water vapour column values, provided by AERONET, were estimated using Sun direct measurements. All mentioned data provided by AERONET can be downloaded visiting the website <http://aeronet.gsfc.nasa.gov>.

The daily total ozone column, TOC, used in this work was measured by several satellite-based remote sensing instruments: OMI collected in the Aura satellite (downloaded from <http://disc.sci.gsfc.nasa.gov/giovanni>) and GOME and GOME-2 data (supplied by the staff of the Remote Sensing Technology Institute, IMF, of the German Aerospace Centre, DLR). These satellite data show, in general, an excellent agreement with reference ground-based measurements (Antón et al., 2009a,b).

3 Modelling cloudless sky radiance

The methodology used in this work for the absolute calibration in radiance of the All-sky Imager is based on the comparison of its output measurements with the sky radiance values estimated by the UVSPEC radiative transfer model under cloud-free conditions. This Section presents the description of the inputs used in the code and the validation of the modelled radiance values against experimental data measured by the CIMEL sunphotometer.

3.1 UVSPEC model

The reference values of sky radiance were estimated by the LibRadtran software package (the 1.6 beta version was used in this paper) which main tool is the UVSPEC model, developed by Mayer and Kylling (2005). The various radiative transfer equation solvers included in the UVPSEC have different capabilities to calculate the radiative quantities in the Earth's atmosphere. The radiative solver chosen to obtain

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sky radiance, on the Earth's surface, was an improved version of the DISORT algorithm in C language (CDISORT) running in 18 stream mode. This solver was developed by Buras et al. (2011), and it uses the correction developed by Nakajima and Tanaka (1988) applying the directional distribution after secondary scattering of light in an atmosphere.

We implemented the UVSPEC model using the standard profiles and inputs indicated by Antón et al. (2011b) for the GFAT station at Granada. The standard atmosphere was mid-latitudes summer (from May to October) and winter (the rest of months). The extraterrestrial spectrum used was the proposed by Gueymard (2005). Surface albedo was considered as a fixed value of 0.2 for all wavelengths. Daily TOC was included in the inputs, and all simulations were run under cloud-free conditions.

For aerosols, the appropriated spring-summer and fall-winter profiles given by Shettle (1989) were used. We chose an urban aerosol profile in the lower 2 km of the atmosphere and a background profile above 2 km, and the impact of the aerosol loading in the boundary was expressed by columnar α and β Angstrom's coefficients (see Sect. 2.3) which are scaled in the chosen profiles. Legendre moments were used as phase function information input, and these moments were calculated using the PMOM tool, included in LibRadtran package, which calculates the Legendre moments of a given phase function. The PMOM tool was run to obtain 200 moments with a scattering angle grid resolution for moderate forward peaks. The SSA and PF as inputs depended on the wavelength of the estimated radiance.

Finally, the UVSPEC outputs for each simulation were the radiances, in $\text{mWm}^{-2}\text{nm}^{-1}\text{sr}^{-1}$, at a selected wavelength each 5° and 1° in azimuth and zenith angles, respectively.

3.2 Modelled vs. experimental radiance

Sky radiances measured by the CIMEL sunphotometer were used to study the reliability of the UVSPEC modelled radiances. For this goal, 50 almucantar cloudless measurements of radiance at 441 and 677 nm were chosen randomly. The cloudless

radiances were calculated using UVSPEC and the inputs explained above, but changing the output angles for the almucantar angles and selecting the wavelengths of the sunphotometer. In addition, SSA441 and PF441 were used as input for the calculations of radiance at 441 nm and SSA677 and PF677 for the radiance at 677 nm.

5 The absolute value of the relative error, ARE, between modelled and experimental values was calculated, for each azimuth angle, by the next equation:

$$\Theta(\%) = 100 \% \frac{|R_{\text{MOD}} - R_{\text{MEAS}}|}{R_{\text{MEAS}}}, \quad (3)$$

where Θ is the ARE, and R_{MOD} and R_{MEAS} are the modelled (UVSPEC) and measured (CIMEL) sky radiances, respectively. ARE values were calculated for each azimuth angle (angle relative to sun azimuth). Figure 3 shows the mean (panel a) and the median (panel b) of ARE using the 50 random values (random solar zenith angles) for the two wavelengths. It can be seen that the mean values of ARE for the sky radiance at 677 nm show values higher than 20% when the azimuth angle is around 0°. In contrast, the behaviour is remarkably good for azimuth angles far from the sun. Thus, the mean parameter varies between 4% and 11% for azimuth angles larger than 12°. Figure 3b shows that the median values of ARE at 677 nm are smaller than 10% for all azimuth angles. This high difference between the median and mean of ARE at 677 nm for cases close to the sun indicates that only a few modelled cases significantly differ with respect to the experimental data when the azimuth angle is around 0°. These particular cases could be associated with not completely cloud-free conditions around the sun. On the other hand, the mean and median values of ARE at 441 nm range are smaller than 15% for all azimuth angles.

Therefore, the sky radiance estimated with the UVSPEC model, and used in this work as reference for the calibration of the All-sky Imager is in agreement with experimental measurements for all azimuth angles. Nevertheless, we have decided to avoid the use of the UVSPEC simulations for those cases near to the sun (azimuth angle smaller than 10°).

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4 Calibration method

This Section describes the method followed to obtain the calibration matrix of the All-Sky Imager, and thus, to convert the output signal of each pixel and channel in the sky radiance that reaches every pixel under cloud-free conditions.

4.1 Effective wavelength

In order to obtain the radiance reaching the All-sky Imager at every pixel, it is necessary to simulate spectral radiance between 400 and 700 nm and weight them using the spectral response shown in Fig. 1 for each channel. This process expends a lot of computation time, moreover, we are interested in obtain spectral radiance instead broadband. This issue is solved using the concept of effective wavelength (Kholopov, 1975). The ratio of two broadband measurements, taken with the same instrument, with its self spectral response, under different conditions, is also equal to the ratio of the same measurements but measured with an instrument which only is sensitive at the effective wavelength, λ_e . Therefore, the effective wavelength for each channel is calculated using the following expression:

$$\lambda_e = \frac{\int \lambda I(\lambda) S(\lambda) d\lambda}{\int I(\lambda) S(\lambda) d\lambda}, \quad (4)$$

where λ is the wavelength, I is the irradiance reaching the instrument, and S is the spectral response of the channel (Fig. 1). In order to calculate λ_e , a different set of irradiances reaching the camera is simulated using the UVSPEC model. A total of 200 simulations of spectral diffuse irradiance (direct irradiance does not reach the camera) were simulated per channel. The phase function used in these simulations is the Henyey-Greenstein function. In addition, the values of the asymmetry parameter and SSA are 0.7 and 0.9, respectively, and a fixed TOC value of 300 DU is utilized for all wavelengths under a mid latitudes summer atmosphere.

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The 200 simulations for each channel (600 in total) were run changing SZA (from 10° to 80° in 10° steps), α parameter (from 0.2 to 1.8 in 0.4 steps) and β parameter (from 0.01 to 0.21 in 0.05 steps). Thus, 200 effective wavelengths per channel were calculated by Eq. (4) using the simulated spectral diffuse irradiances. The mean values were 464 nm (blue), 534 nm (green) and 626 nm (red), with a standard deviation of 2 nm for the three channels. Therefore, the spectral radiance reaching the camera is simulated at these three effective wavelengths as is described in Sect. 3.1, taking into account that the SSA and PF values used as inputs are SSA441 and PF441 to estimate the radiance at 464 nm, SSA550 and PF550 for 534 nm, and SSA677 and PF677 for 626 nm.

4.2 Calibration matrix

First, a cloudless image is selected, being separated in the red, green and blue images. We mask the solar zenith angles higher than 80° , the pixels near to the sun, and the different obstacles around the whole sky like the shadow system and the two pyranometers installed on the tracking system near the camera (see Fig. 2). In the next step, the pixel counts were normalized to unity using the ratio between the raw value and the highest recorded value (2^{16}). Finally, the dark noise signal was removed taking into account that, in all rows and columns of an image, a dark zone appears whose signal must be null (as a first approximation). Therefore, the minimum values for each column and each row were averaged and considered as the dark signal, which was subtracted to the normalized signal. Some authors (e.g. Voss and Zibordi, 1988; Voss and Liu, 1997) found problems related to the lens since its transmittance varies with the field of view. However, in this work, this issue was not considered because the calibrated system consists of the CCD with the lens together, and the field of view of the lens presents no significant changes in the pixels between different images due to the angular symmetry.

The corrected image is a 900×900 matrix, \mathbf{P} , which P_{ij} element is the corrected raw signal for the ij -pixel. This signal should be proportional to the incident radiation:

radiance that reaches the ij -pixel multiplied by the solid angle that view, Ω_{ij} . Therefore, a 900×900 matrix, \mathbf{R} , was constructed using the simulated radiance values (Sects. 3.1 and 4.1) and interpolating for each pixel. Therefore, the matrix calibration, \mathbf{K} , can be expressed as:

$$\mathbf{K} = \frac{\mathbf{R} \Omega}{\mathbf{P}}. \quad (5)$$

This relationship is only valid if the response of the CCD sensor is lineal. The raw signal of the CCD, without removing the dark noise, was represented as a function of the simulated sky radiance for different cloudless images (not shown). The results indicate that the CCD response is linear for normalized raw values smaller than 0.8, being the pixel saturated for a signal higher than 0.8. Therefore, the pixels with normalized signal higher than 0.8 (before the subtraction of the dark noise) were removed along with their eight neighbours due to blooming effect (Voss and Liu, 1997).

From the method described here, given a cloudless image, three calibration matrices (\mathbf{K} -RED, \mathbf{K} -GREEN and \mathbf{K} -BLUE) are calculated using Eq. (5). These matrices can be used to calculate the radiance at 464, 534 and 626 nm multiplying the specific \mathbf{K} matrix by the specific channel of the image and dividing this value by the solid angle viewed for each pixel.

4.3 Matrix calibration variability

A set of cloudless images were selected to study the variability of the matrix elements (K_{ij}). The dates of the eight chosen days along with the initial and final time of their images set (every 5 min) are presented in Table 1. These days were selected to study the variability between consecutive days (day 1 to 4), and the variability in a long period (5 months) from winter days (day 1 to 6) to summer ones (day 7 and 8). Moreover, the eight selected days were completely cloud-free in order to see the intra-daily variation of the calibration matrix. Therefore, a \mathbf{K} matrix can be obtained for each image recorded during the eight days shown in Table 1.

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In order to analyze the influence of SZA on the K_{ij} elements, we evaluated the variability for each ij -pixel over a whole day. Thus, the intra-daily coefficient of variation, Γ , was calculated for each ij -pixel as:

$$\Gamma_{ij} = 100 \% \frac{\sigma(K_{ij,T})}{M(K_{ij,T})}, \quad (6)$$

5 where $\sigma(K_{ij,T})$ and $M(K_{ij,T})$ are the standard deviation and the mean of the several values of K_{ij} throughout the day. Thus, Γ is a matrix whose elements indicate the daily variation of each K_{ij} .

Table 2 shows the percentage of K_{ij} elements with Γ lower than 10 % and the mean value of Γ for each day. It can be seen that the intra-daily variation of the coefficients of the matrix calibration is small, with a mean value between 3.1 % and 5.4 % for all days and channels. In addition, the percentage of coefficients with variations lower than 10 % is near to 100 % for all days. These results justify that we assume that **K-RED**, **K-GREEN** and **K-BLUE** does not depend on SZA, and that daily mean **K** matrix, \mathbf{K}_d , obtained as the average of all calculated **K** matrices in a day can be considered
 15 representative for that day. Thus, \mathbf{K}_d is calculated for the eight selected days, and the resulting matrices are compared to each other. The matrix with the absolute difference between the \mathbf{K}_d of the m -day ($\mathbf{K}_{d,m}$) and the n -day ($\mathbf{K}_{d,n}$) is obtained as:

$$\Delta\mathbf{K}_{d,m,n}(\%) = 100 \% \frac{|\mathbf{K}_{d,m} - \mathbf{K}_{d,n}|}{\mathbf{K}_{d,n}}. \quad (7)$$

Table 3 reports about the inter-comparison between the eight days, showing the percentage of pixels with a difference lower than 10 % ($\Delta\mathbf{K}_{d,m,n} < 10\%$) and the mean of these differences in parentheses. The mean difference between the days 1, 2 and 3 are smaller than 2 % with a 99.9 % of pixels with a difference lower than 10 %. This result indicates low variation in **K** between consecutive days. The highest mean difference between two winter days is 2.9 %, 4.0 % and 4.5 % for the red, green and blue channels, respectively. Moreover, the percentage of pixel with a difference lower than
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534 and 626 nm. To solve these differences in the wavelengths of the channels, we obtain the ratios between the radiances at 677 nm and 626 nm (Ratio R), 501 nm and 534 nm (Ratio G), and 441 nm and 464 nm (Ratio B) using the same 200 spectra of diffuse irradiance calculated in Sect. 4.1. The averages (\pm standard deviation) of these ratios are 0.87 ± 0.07 (Ratio R), 1.13 ± 0.06 (Ratio G) and 0.98 ± 0.05 (Ratio B). The camera radiance in a given solar direction was obtained as the average of the closest pixels in that direction (a square of 25 pixels). Thus, the multiplication of these radiances for the three averages ratios results in the estimation of the camera radiance at the CIMEL wavelengths.

Figure 5a–f shows the cloudless sky radiance measured by the sunphotometer and the estimations given by camera for the almucantar (left) and principal plane (right). Each panel corresponds to a particular case which was randomly selected with the unique condition that the time difference between the CIMEL and the camera measurements must be smaller than 10 min. It can be seen that CIMEL and camera radiances show a similar behaviour. To evaluate the camera-CIMEL differences, the ARE mean values are calculated using 40 random cloudless images for 677 nm and 441 nm and 11 images for 501 nm. Figure 5 shows the variability of this parameter for the almucantar (panel g) and principal plane (panel h). It can be observed that the errors are lower than 15% for all azimuth angles except near to the sun when the disagreement experiences a notable increase. For the three wavelengths, the errors in the principal plane are also smaller than 15% except for the lowest and highest zenith angles. These errors are similar than the obtained between the UVSPEC simulations and CIMEL (Fig. 3), therefore, the proposed method and the obtained \mathbf{K} matrices can be consider valid to estimate sky radiances.

5.2 Application of the calibration matrices

Spectral sky radiances were calculated for three cases shown in Fig. 2, which correspond to three different sky conditions.

(± 30), 60 (± 30) and 90 (± 40) $\text{m Wm}^{-2} \text{nm}^{-1} \text{sr}^{-1}$ for the red, green and blue channel, respectively. The highest deviation is caused by the differences between cloudless and cloud cover pixels. The blue channel shows a similar mean value than in the cloudless and overcast conditions, which is due to its wavelength is more affected by molecular scattering and, therefore, diffuse radiation is high also under cloudless conditions. However clouds strongly increase the higher wavelengths.

6 Conclusions

Some important conclusions may be drawn from this work. The radiative transfer model UVSPEC estimates radiance values, under cloudless sky, similar to the recorded with a sunphotometer, obtaining the worse agreement between modelled and measured radiances near to the sun. The mean differences between the modelled and measured sky radiances are lower than 15 % for all azimuth angles when the median is considered. Thus, UVSPEC model can be used to estimate cloudless sky radiances if some inputs, such as aerosol scattering phase function, are taken into account.

The obtained calibration matrices, which convert raw signal in physical radiance, show no significant dependence on SZA, and the variability of their coefficients along 5 months is low, being the coefficient of variation lower than 10 % for the 99 % of the pixels. The mean differences between the radiance estimated with the camera and the measured with the CIMEL are lower than 15 % except for high zenith angles and close to the Sun, showing the radiance at 464 nm the lowest differences in these angles.

The radiance under overcast conditions presents the highest homogeneity, while the largest variability in the radiance values corresponds to the partially cloudy conditions. Individual clouds increase the sky radiance at the higher wavelengths, and the horizon brightening under cloudless conditions changes to horizon dimming when the sky is obscured overcast.

Finally, in future works, other analysis of sky radiances might be developed, e.g. trying to retrieval aerosol properties using camera radiance information. The authors

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encourage researchers to apply and use the proposed method to obtain sky radiance from sky images.

Acknowledgements. The authors gratefully acknowledge the financial support extended by the Spanish Innovation and Science Ministry for the projects: CGL2011-25363 and CGL2010-12140E. Roberto Román thanks Valladolid University for the support to research short stays. This work was also partially supported by the Andalusian Regional Government through projects P08-RNM-3568 and P10-RNM-6299, the Spanish Ministry of Science and Technology through projects CGL2010-18782 and CSD2007-00067, and by the European Union through ACTRIS project (EU INFRA-2010-1.1.16–262254). The Remote Sensing Technology Institute (IMF) at the German Aerospace Centre (DLR) of GOME and GOME-2 instruments and the National Aeronautics and Space Administration (NASA) of OMI are also acknowledged for making available the ozone data required for this study.

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Table 2. Percentage of elements of Γ lower than 10 %, for each day and channel. The mean value of Γ is given in parentheses.

Day	Red	Green	Blue
1	97.1 (5.4)	99.1 (4.4)	100 (3.9)
2	99.1 (5.3)	99.4 (4.3)	99.3 (4.1)
3	99.9 (4.5)	100 (3.5)	100 (3.1)
4	99.8 (4.9)	99.9 (3.9)	99.9 (3.5)
5	100 (4.8)	100 (4.1)	100 (4.5)
6	99.0 (5.3)	100 (4.2)	99.9 (3.9)
7	98.0 (5.3)	99.2 (4.9)	96.0 (5.2)
8	99.6 (4.0)	100 (3.4)	100 (3.5)

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Table 3. Contingency table showing the percentage of elements of $\Delta K_{d,m,n}$ smaller than 10% as a function of days and channels. The mean $\Delta K_{d,m,n}$ is given in the parentheses.

Channel	Day	1	2	3	4	5	6	7	8
Red	1	–	100 (1.3)	100 (1.3)	99.9 (2.9)	99.7 (2.0)	99.6 (1.9)	78.5 (6.1)	99.0 (2.9)
	2	100 (1.3)	–	100 (1.0)	99.8 (2.0)	99.4 (1.8)	99.6 (1.3)	88.9 (5.3)	98.5 (2.6)
	3	100 (1.3)	100 (1.0)	–	99.8 (2.0)	99.1 (2.3)	99.6 (1.6)	86.2 (5.5)	98.9 (2.7)
	4	99.9 (2.8)	100 (2.0)	99.9 (1.9)	–	99.5 (3.0)	99.6 (2.0)	93.6 (4.9)	97.5 (3.5)
	5	99.6 (2.0)	99.4 (1.8)	99.5 (2.3)	99.4 (3.1)	–	99.3 (1.7)	87.0 (5.5)	98.6 (2.7)
	6	99.6 (1.9)	99.6 (1.3)	99.6 (1.5)	99.6 (2.1)	99.4 (1.6)	–	91.6 (5.0)	98.0 (2.5)
	7	83.2 (5.8)	93.5 (5.1)	91.3 (5.3)	95.5 (4.9)	92.0 (5.3)	95.7 (4.8)	–	98.0 (3.8)
	8	98.7 (2.9)	98.1 (2.7)	98.7 (2.7)	97.1 (3.7)	97.9 (2.7)	97.6 (2.6)	96.6 (4.0)	–
Green	1	–	100 (1.4)	100 (1.9)	99.9 (4.0)	99.9 (1.5)	99.9 (1.7)	69.1 (7.0)	99.5 (4.1)
	2	100 (1.4)	–	99.9 (0.8)	99.9 (2.6)	99.9 (1.9)	99.8 (1.0)	77.8 (6.1)	99.9 (3.2)
	3	99.9 (1.8)	99.9 (0.8)	–	99.8 (2.2)	99.9 (2.3)	99.8 (1.2)	79.4 (5.9)	99.7 (3.1)
	4	100 (3.8)	99.9 (2.5)	99.8 (2.2)	–	99.8 (4.0)	99.7 (2.8)	91.6 (5.0)	99.8 (2.6)
	5	99.9 (1.5)	99.8 (1.9)	99.9 (2.4)	99.7 (4.2)	–	100 (1.6)	70.7 (6.6)	99.9 (3.9)
	6	99.8 (1.6)	99.8 (1.0)	99.8 (1.2)	99.6 (2.9)	100 (1.6)	–	78.7 (5.8)	99.8 (2.9)
	7	73.8 (6.4)	83.2 (5.7)	85.9 (5.6)	95.7 (4.9)	77.5 (6.0)	86.1 (5.4)	–	99.5 (3.0)
	8	99.9 (3.9)	99.9 (3.1)	99.8 (3.0)	99.8 (2.7)	100 (3.7)	99.8 (2.8)	98.9 (3.2)	–
Blue	1	–	100 (1.3)	100 (2.0)	99.9 (4.5)	99.8 (1.4)	99.8 (1.6)	73.6 (6.5)	94.0 (4.8)
	2	100 (1.3)	–	100 (0.8)	99.9 (3.3)	99.8 (1.8)	99.7 (1.1)	80.9 (5.8)	99.6 (3.8)
	3	100 (2.0)	100 (0.8)	–	99.8 (2.7)	99.8 (2.0)	99.7 (1.3)	85.4 (5.4)	99.6 (3.4)
	4	99.9 (4.3)	99.9 (3.2)	99.9 (2.6)	–	99.4 (4.3)	99.7 (3.4)	94.5 (4.9)	99.7 (2.8)
	5	99.8 (1.4)	99.7 (1.8)	99.7 (2.1)	98.6 (4.5)	–	100 (1.5)	78.8 (5.8)	97.7 (4.5)
	6	99.8 (1.6)	99.7 (1.1)	99.7 (1.3)	99.4 (3.5)	100 (1.5)	–	84.8 (5.3)	99.8 (3.6)
	7	78.0 (6.0)	87.2 (5.4)	90.9 (5.1)	95.6 (4.9)	85.0 (5.4)	90.2 (5.0)	–	99.9 (2.6)
	8	98.4 (4.5)	99.8 (3.6)	99.8 (3.3)	99.6 (2.8)	99.6 (4.3)	99.8 (3.4)	99.8 (2.6)	–

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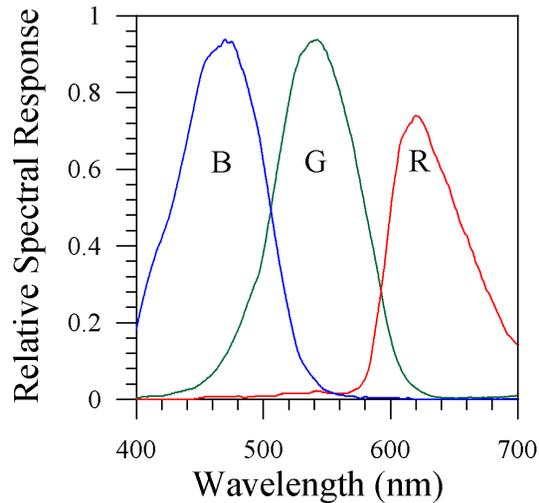


Fig. 1. Relative spectral response of the ICX085AK CCD sensor for the red (R), green (G) and blue (B) channels, taking into account the infrared filter included in the RETIGA 1300C camera.

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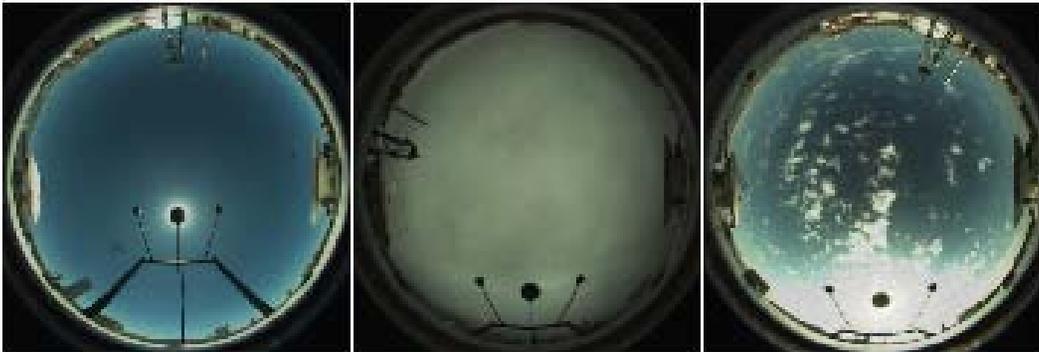


Fig. 2. Sky images taken by the All-sky Imager of GFAT for three different conditions: a cloudless sky (8 July 2011, 13:15 UTC) on the left, an obscured overcast sky (23 January 2011, 10:35 UTC) in the middle, and a partially cloudy sky (7 March 2011, 16:15 UTC) on the right.

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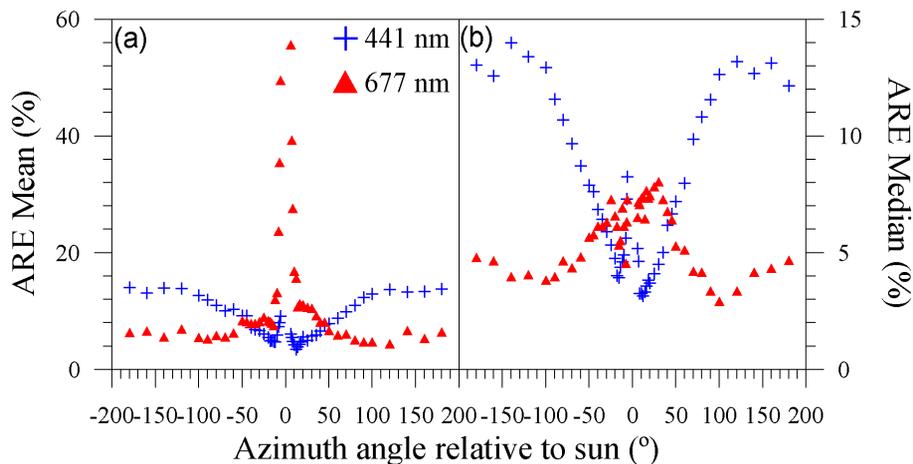


Fig. 3. Mean (a) and median (b) of the absolute relative error for 441 and 677 nm as a function of the azimuth angle relative to sun.

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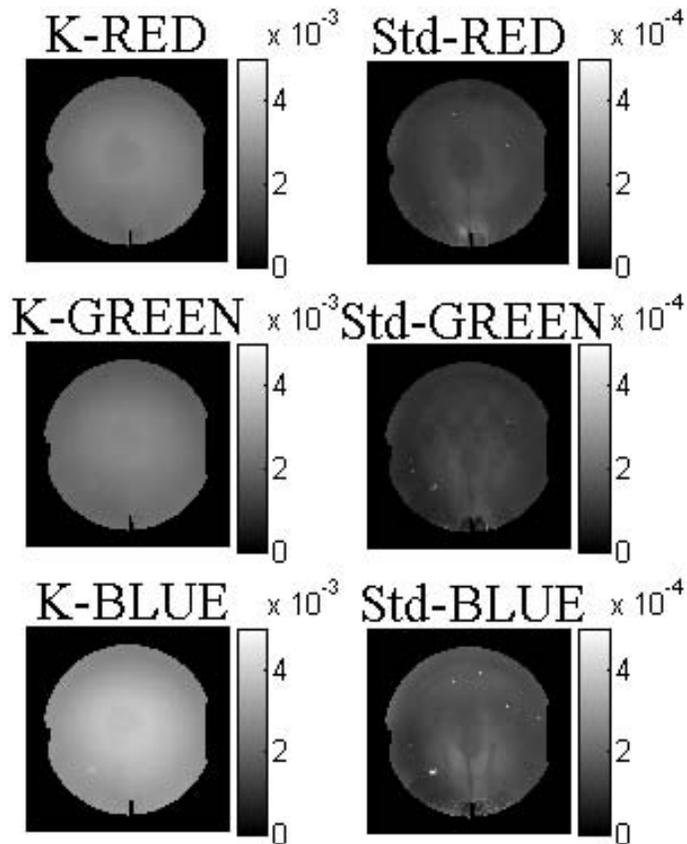


Fig. 4. Calibration matrices and their standard deviation (Std) in $\text{m Wm}^{-2} \text{nm}^{-1} \text{raw}^{-1}$, for the three channels.

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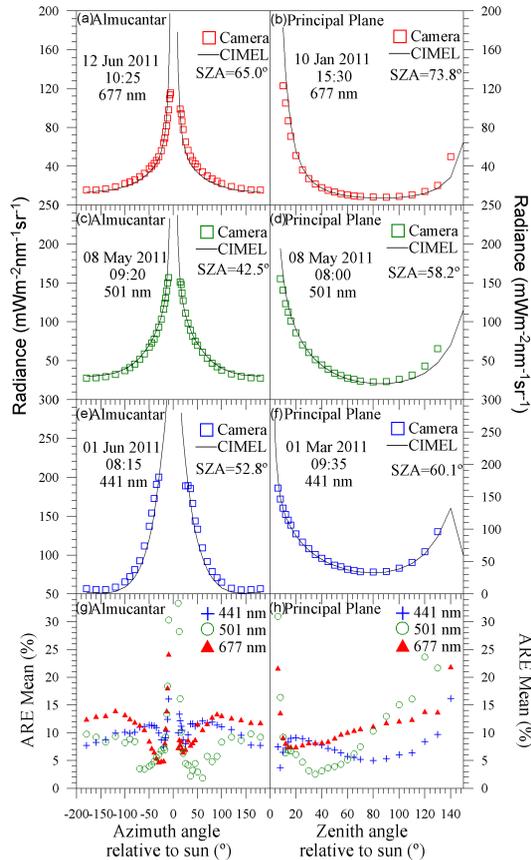


Fig. 5. The CIMEL and camera sky radiances together for two different dates at 677 nm (a, b), 501 nm (c, d) and 441 nm (e, f). (g) and (h) shows the mean of the absolute relative error for 677 nm, 501 nm, and 441 nm as a function of the azimuth and zenith angle. Left panels represent almucantar and right panels are principal planes. The zenith angle equal to zero is the SZA, and the highest zenith angles correspond to the azimuth of 180°.

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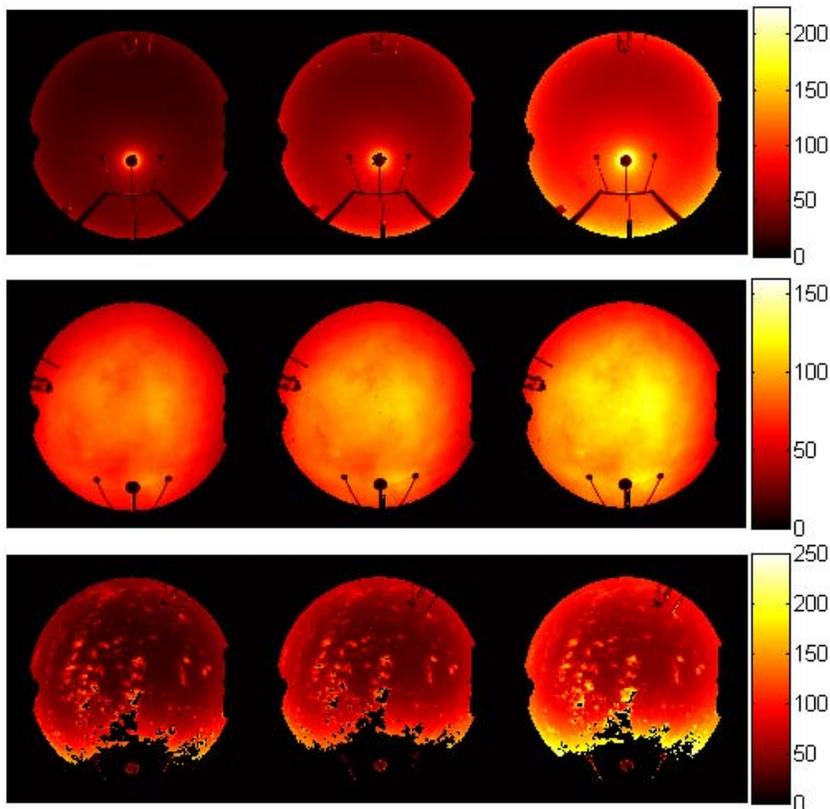


Fig. 6. Sky radiance (in $\text{m Wm}^{-2} \text{ nm}^{-1} \text{ sr}^{-1}$) at 626 nm (left panel), 534 nm (centre panel) and 464 nm (right panel) for Top panel: a cloudless sky (8 July 2011 at 13:15 UTC), Middle panel: an obscured overcast sky (23 January 2011, 10:35 UTC), and Bottom: a partially cloudy sky (7 March 2011, 16:15 UTC). Black regions represent the zenith angles higher than 80° , the two pyranometers installed near the camera, and the saturated pixels. Figure 2 helps to discern black saturated (white) and low-radiance (black) pixels.

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